

# Single frequency blue laser fiber amplifier

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**An all-fiber amplifier for single-frequency blue laser was demonstrated for the first time to the best of our knowledge. Over 150 mW continuous-wave single-transverse-mode blue laser output was obtained with a 10-m 1000 ppm thulium-doped fluoride fiber pumped by a 1125 nm fiber laser at a power of 2 W. The output power was limited due to the onset of the competitive lasing at 783 nm. Photo-darkening and photo-curing of the thulium-doped fiber amplifier were also studied and analyzed. © 2017 Optical Society of America**

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Thulium (Tm<sup>3+</sup>) doped ZBLAN (ZrF<sub>4</sub>-BaF<sub>2</sub>-LaF<sub>3</sub>-AlF<sub>3</sub>-NaF) features a large number of possible laser transitions with emission wavelengths ranging from ultraviolet (UV) to mid-infrared due to the low phonon energy of ZBLAN (550 cm<sup>-1</sup>), while these transitions are typically non-radiative in Tm<sup>3+</sup> doped silica [1]. For instance, the transition from state <sup>3</sup>H<sub>4</sub> to state <sup>3</sup>F<sub>4</sub> can generate laser at 1480 nm that can be used for optical communications, while the transition from state <sup>3</sup>H<sub>4</sub> to state <sup>3</sup>H<sub>5</sub> can produce light around 2.3 μm making it a good candidate for hydrocarbon gases detection. In addition, transitions between <sup>1</sup>G<sub>4</sub> and <sup>3</sup>H<sub>5</sub> states and <sup>3</sup>H<sub>4</sub> and <sup>3</sup>H<sub>6</sub> states have broad emission in a band around 800 nm which is in demand for medical applications. However, the most attractive property of Tm<sup>3+</sup>-doped ZBLAN is its capability of lasing at UV and blue wavelengths [1], where robust and reliable laser sources are needed for a variety of applications including light detection and ranging (LIDAR) [2], optical data storage [3], imaging [4], and spectroscopy. Tm<sup>3+</sup>-doped ZBLAN fiber lasers have been extensively studied in 1990s for making compact blue lasers for these applications considering the advantages of fiber

lasers such as guided laser beam, high power scalability, very low thermal effects, low maintenance, and compactness.

The first Tm<sup>3+</sup>-ZBLAN fiber laser in the blue was demonstrated by Grubb *et al* [5]. A Nd<sup>3+</sup>:YAG laser at 1120 nm was used as the pump source and 60 mW output was obtained at 480 nm. After that, various Tm<sup>3+</sup>-ZBLAN fiber lasers were demonstrated with either increased output power, or increased efficiency, or reduced threshold [1]. However, to the best of our knowledge, the lasers of all these experiments were neither single transverse mode nor single longitudinal mode. Single frequency blue lasers are in great demand for optical metrology, interferometry, high-order harmonic generation, quantum cryptography, high resolution spectroscopy, and under-sea communication. In this paper we report the first demonstration of an all-fiber blue laser amplifier with single frequency and single transverse mode output.

The energy level diagram of Tm<sup>3+</sup> is shown in Fig. 1(a). Blue laser emission around 480 nm can be obtained by directly exciting the ground state <sup>3</sup>H<sub>6</sub> to <sup>1</sup>G<sub>4</sub> by 465 nm pumping. Since there is no readily available high power single-mode laser source at 465 nm, this is not a practical option at this time. Therefore, upconversion pumping at 1125 nm was used in our experiment. Due to the long lifetimes of some metastable energy levels between <sup>3</sup>H<sub>6</sub> and <sup>1</sup>G<sub>4</sub>, ions can be excited to the upper laser level <sup>1</sup>G<sub>4</sub> by absorbing three pump photons as shown in Fig. 1(a). There are three radiative transitions from energy level <sup>1</sup>G<sub>4</sub> to lower energy levels with emission wavelengths at 480 nm, 650 nm and 784 nm. The branching ratios of the three transitions are 0.378, 0.077, and 0.384, respectively [6]. Fig. 1(b) shows the measured fluorescence from a Tm<sup>3+</sup>-doped ZBLAN glass. The emissions at 480 nm and 650 nm corresponds to the transitions from state <sup>1</sup>G<sub>4</sub> to <sup>3</sup>H<sub>6</sub> and <sup>3</sup>F<sub>4</sub>, respectively. The emission peak near 800 nm, however, is the combination of two transitions: <sup>1</sup>G<sub>4</sub> → <sup>3</sup>H<sub>5</sub> and <sup>3</sup>H<sub>4</sub> → <sup>3</sup>H<sub>6</sub>. Both transitions influence the operation of the 480 nm amplifier. Population inversion for the 784 nm transition is much easier achieved than that for the 480 nm emission due to the empty

lower level  $^3H_5$ . On the other hand, 810 nm emission depletes the population of the  $^3H_4$  state that is important for the multiphoton excitation. Therefore, a technique has to be implemented in order to suppress the competitive lasing at the 800 nm band.

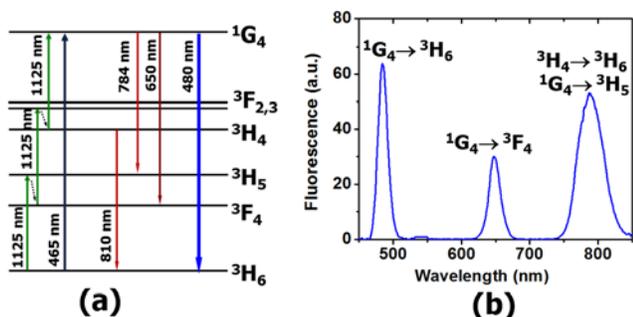


Fig. 1. (a) Partial energy level diagram of  $Tm^{3+}$  and transitions related to the blue laser emission; (b) Measured fluorescence of  $Tm^{3+}$ -ZBLAN glass.

In order to obtain single transverse mode output, a  $Tm^{3+}$ -doped ZBLAN fiber with a core diameter of 4  $\mu m$  and a numerical aperture of 0.07, corresponding to a V-number and cutoff wavelength of 2.095 and 365 nm respectively, was designed and fabricated. The microscope image of the end facet of the fiber is shown in Fig. 2(a). This fiber has a 12- $\mu m$  inner cladding doped with 100 ppm neodymium ( $Nd^{3+}$ ) which intentionally suppresses 800 nm emission due to  $Nd^{3+}$  absorption in this wavelength region. The outer cladding of this fiber is 125  $\mu m$  compatible with standard silica fiber enabling easy fusion splicing and the construction of all-fiber amplifiers. Fig. 2 (b) shows the 2D diagram of refractive index profile of the fiber measured with an interferometric fiber analyzer (Interfiber Analysis, IFA-100). The refractive index difference between core and cladding is about 0.002. The fiber core is doped with 1000 ppm  $Tm^{3+}$  and the absorption of this fiber of 1125 nm light launched into the fiber core was measured at a low power level by a cutback experiment to be about 1.56 dB/m.

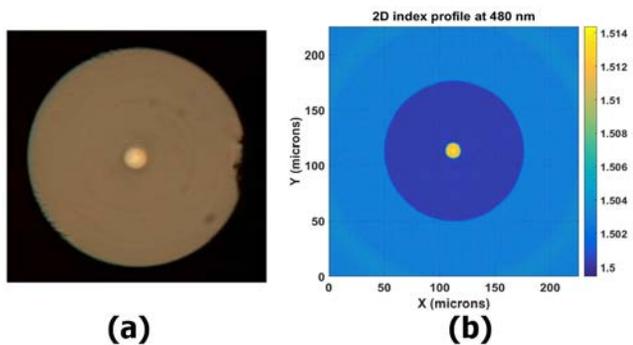


Fig. 2. (a) Microscope image of an end face and (b) 2D refractive index profile of the single-mode  $Tm^{3+}$ -doped ZBLAN fiber.

The schematic of the all-fiber blue laser amplifier is shown in Fig. 3. A single frequency external cavity GaN diode laser at 478 nm coupled to a polarization maintaining (PM) single-mode silica fiber (PM 460HP) was used as the seed laser. An ytterbium ( $Yb^{3+}$ ) doped silica fiber laser operating at 1125 nm was used as the pump source. A PM wavelength division multiplexor (WDM) made of PM 460HP (signal port) and PM980 fiber (pump and common port) was used to combine the seed signal laser and the pump laser together. The signal coupling efficiency is about 80% and the pump coupling efficiency is about 92%. The polarization extinction ratio (PER) of the signal laser changes from an initial value of 25 dB to 18 dB after the PM WDM indicating that most signal laser is coupled to the fundamental mode of the PM980 fiber although this fiber has a cutoff wavelength of 705 nm and is multimode for the blue light. It should be noted that the common port fiber of the PM-WDM was fixed on the optical table to reduce power fluctuations and decrease of the fundamental mode caused by small bending and other environmental effects. 10-m of  $Tm^{3+}$ -ZBLAN fiber was spliced to the common port PM980 fiber using NP Photonics proprietary splicing technique [7]. The image of the splice joint between the PM980 silica fiber and the  $Tm^{3+}$ -doped ZBLAN fiber is shown in the inset of Fig. 3. The typical loss of the splice between a ZBLAN fiber and a silica fiber is less than 0.3 dB. 10-m ZBLAN fiber was chosen based on the cutback experiment because it produced the maximum output power. The output end of the ZBLAN fiber was angle cleaved to reduce backward reflections. A band-pass filter at 480 nm was used to remove the residual pump at 1125 nm.

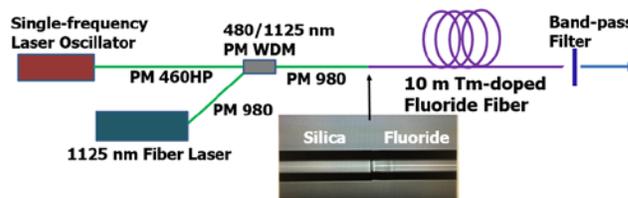


Fig. 3. Experimental setup of the all-fiber blue laser amplifier. Inset: the picture of the splice joint between the silica fiber and the fluoride fiber.

In the experiment, single-frequency operation of the GaN diode laser was confirmed by a scanning Fabry-Perot interferometer (FPI) as shown in the lower right inset of Fig. 4. The blue signal laser power after the PM-WDM was approximately 15 mW. The optical spectra of the fiber amplifier with and without pump were measured with an optical spectrum analyzer (Ando, AQ6351A) and is shown in the upper left inset of Fig. 4. The output power of the fiber amplifier as a function of the pump power was measured by a power meter (Thorlabs, PM100D) and is shown in Fig. 4. The pump threshold of the fiber amplifier is about 750 mW. The blue laser output increases with pump power and saturates when the pump power exceeds 2 W. A maximum output power of 155 mW, corresponding to a net gain of 10 dB was obtained. The slope efficiency of this fiber amplifier is about 13%, which is smaller than those of previous experiments because of smaller overlaps of pump and signal lasers and the lower fiber NA when compared to multimode fibers used in previous reports. The lower efficiency is the result of the tradeoffs employed to achieve single transverse mode output beam using a low core NA fiber.

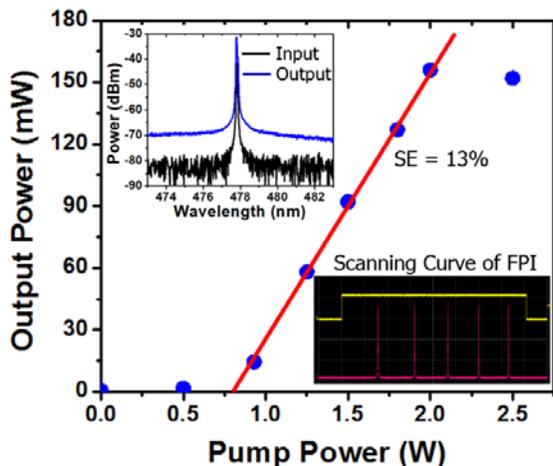


Fig. 4. Measured output power of the 478 nm fiber amplifier as a function of the launched pump power. Upper left inset: Optical spectra of the input and output of the fiber amplifier. Lower right inset: Single-frequency operation of the GaN diode laser confirmed with a scanning Fabry-Perot Interferometer.

The beam quality of the fiber amplifier output was measured with a beam profiler (DataRay Inc., Beam Map2). The 2D beam profile is shown in Fig. 5(a). The cross-section of the beam profile and its Gaussian fit are shown in Fig. 5(b). There are two small side peaks in the measurement due to excited cladding modes. It should be noted that the beam profile was not stable and the measured beam quality was about  $M2 \sim 1.5$  at low pump powers because of interference between the cladding modes and the fundamental core mode supported by the long coherence length of the single-frequency signal laser with a typical linewidth  $< 1$  MHz. The beam quality improved with increasing pump power. A beam quality  $M2$  of 1.06 was measured at the maximum output power. The PER of the output laser was measured to be approximately 11 dB even though the  $\text{Tm}^{3+}$ -ZBLAN fiber is not designed to maintain the polarization of the signal laser.

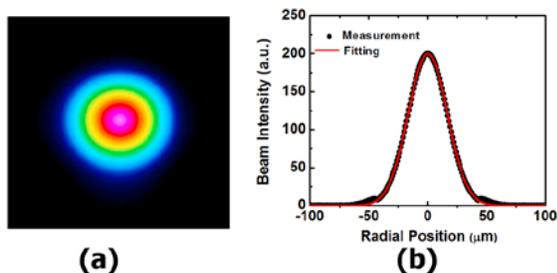


Fig. 5. (a) 2D profile of the fiber amplifier output beam; (b) Cross-section of the output beam profile and its Gaussian fit.

As mentioned above, the 12- $\mu\text{m}$  inner cladding of the  $\text{Tm}^{3+}$ -doped ZBLAN fiber was doped with 100 ppm  $\text{Nd}^{3+}$  in order to suppress the competitive 800 nm laser.  $\text{Nd}^{3+}$  has significant absorption near 800 nm due to transitions from  $^4I_{9/2}$  to  $(^2H_{9/2}, ^4F_{5/2})$  as shown by the black curve in Fig. 6. Absorption of 800 nm emission from  $\text{Tm}^{3+}$  occurs because of partial mode overlap

between the core mode field and the  $\text{Nd}^{3+}$  doped cladding. However, the output power of this  $\text{Tm}^{3+}$ -ZBLAN fiber amplifier still saturated at a pump power of 2 W and decreased with further increasing pump power due to the onset of the laser at 783 nm as shown in Fig. 6. When the pump power is 2 W, there is ASE at 800 nm and  $\text{Tm}^{3+}$  long wavelength lasing is effectively suppressed by the  $\text{Nd}^{3+}$  absorption in the inner cladding. When the pump power is 2.5 W, the 783 nm laser starts and the 478 nm signal power decreases as shown by the blue curve. Further suppressing the 783 nm laser may be obtained with an optimized  $\text{Nd}^{3+}$  doping level but we should notice that  $\text{Nd}^{3+}$  also has some absorption at 478 nm, which reduces the efficiency of the fiber amplifier correspondingly. Proper fiber design that can produce high loss at 783 nm while negligible loss at the wavelength of the blue laser could offer an attractive option of further suppress the competitive lasing.

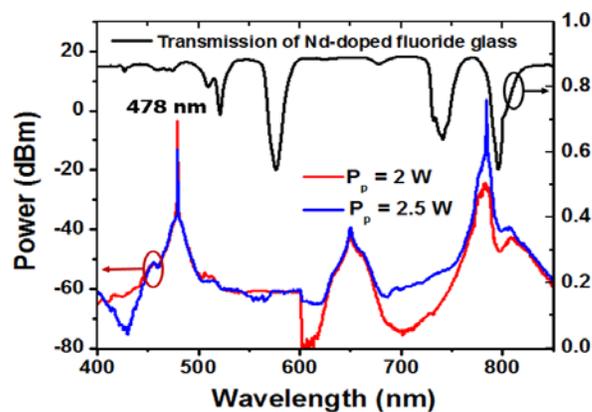


Fig. 6. Optical spectra of the fiber amplifier when the pump power is 2 W (red curve) and 2.5 W (blue curve), respectively. The black curve: Transmission of a 1-cm thick  $\text{Tm}^{3+}$ -doped ZBLAN glass.

Another obstacle to achieving long-term and stable high power single mode blue laser output is the degradation of the  $\text{Tm}^{3+}$ -doped ZBLAN fiber caused by the photodarkening effect. Photodarkening describes a significant increase of background loss at visible wavelengths due to the creation of color centers in the ZBLAN glass matrix of ZBLAN glass under illumination with a strong blue laser [8, 9]. The photodarkening effect in  $\text{Tm}^{3+}$ -doped ZBLAN is particularly strong because  $\text{Tm}^{3+}$  has a large range of energy levels including levels with energies higher than  $^1G_4$  that are capable of absorbing and emitting UV light. In our experiment, it was found that blue laser signal transmitted through the  $\text{Tm}^{3+}$ -doped ZBLAN fiber amplifier significantly decreased (from a few mW to tens of  $\mu\text{W}$ ) after running at the maximum output power and then leaving unpumped for one day. However, the photodarkened  $\text{Tm}^{3+}$ -doped ZBLAN fiber can be photo-cured by launching the blue laser into the fiber core and pumping it at a low pump power for a few hours. The maximum output power of the photo-cured fiber amplifier could be 75% of the initial maximum output power (155 mW). It was also found that photodarkening and photo-curing of the  $\text{Tm}^{3+}$ -doped ZBLAN fiber were repeatable.

In our experiment, the photodarkening effect in the  $\text{Tm}^{3+}$ -doped ZBLAN fiber was found to be not uniform through the length of the

fiber. The transmission spectra of the first 1-m segment at the beginning and the last 1-m segment at the end of the photodarkened amplifier fiber, and a 1-m of fresh fiber (never pumped) were measured using a white light source and an OSA and are shown in Fig. 7. Contrary to our initial expectation, the first 1-m segment of the amplifier fiber, where the blue laser power is low, was severely photodarkened while the last 1-m segment of the amplifier fiber, where the blue laser power is much higher, experienced little photo-darkening, indicating that the photodarkening is not solely induced by the blue laser and may be a function of the population at the upper laser level  $^1G_4$ . It was also observed that photodarkening occurs only after the parasitic 800 nm laser begins. Several previous papers have tried to help in the understanding of photodarkening [8, 9]. But none of them can be used to explain the phenomena observed in our fiber amplifier experiment. Therefore, further investigation and new interpretations of photodarkening in  $Tm^{3+}$ -doped ZBLAN fiber are essential for further developments of high power fiber amplifier for blue lasers.

In conclusion, an all-fiber amplifier for single-frequency blue laser was demonstrated for the first time. More than 150 mW continuous-wave, single frequency and single-transverse-mode 478 nm laser output was obtained with an 1125 nm pump laser at a power of 2 W. The output power was saturated and began to decrease for the pump powers exceeded 2 W due to the onset of the 800 nm lasing. It was observed that the effect of photodarkening is not uniform across the amplifier fiber and further investigations of photodarkening in  $Tm$ -doped ZBLAN fiber amplifier is essential. Further power scaling of the  $Tm^{3+}$ -ZBLAN fiber amplifier can be achieved by sufficiently suppressing the 800 nm laser with optimized  $Nd^{3+}$  concentration in the inner cladding and a proper fiber design with tailored guide at different wavelengths.

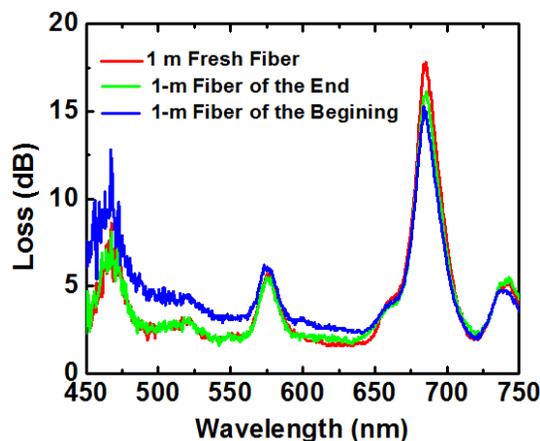


Fig. 7. Measured loss of 1 meter fresh (unpumped) fiber, 1 meter segment of the end and 1 meter segment of the beginning of the photodarkened amplifier fiber.

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